

# Lecture 12: Stochastic integrals

## 1 Introduction

Evans [1] discusses stochastic integral in § B and § C of chapter IV. In § D [1] derives Ito lemma as a result for differentials along paths specified by non-anticipative functionals of the Wiener process. We showed in lecture 9 that Ito lemma can be regarded as a result in standard analysis for functions of paths with *finite second variation*. Chapter IV of [2] in § 4.1–4.2 also covers the construction of stochastic integrals.

## 2 Stochastic integrals

Let  $f$  an analytic function

$$f: \mathbb{R} \rightarrow \mathbb{R}$$

we would like to make sense of the functional of the Brownian motion

$$I = \int_0^t f(w_s) dw_s \quad \text{mathematics notation} \quad (2.1)$$

sometimes also written as

$$I = \int_0^t f(w_s) \eta_s ds \quad \text{physics notation} \quad (2.2)$$

Note that the physics notation implies that  $\eta_s$  is the “derivative” of the Wiener process which we showed to be nowhere differentiable. We should therefore interpret  $\eta_s ds$  only as an alternative notation for  $dw_s$ . It is important to realise that the integral on the right hand side **cannot be interpreted as a Lebesgue-Stieltjes integral**.

### 2.1 Example: Wiener process as integrand

Namely take

$$f(w_s) = w_s$$

and suppose to define the integral as

$$\int_0^t w_s^{(\theta)} dw_s = \lim_{|\mathcal{P}| \downarrow 0} \sum_{t_k \in \mathcal{P}} w_{\theta_k} (w_{t_k} - w_{t_{k-1}}) \quad (2.3)$$

As  $n$  increases the  $\{t_k\}_{k=1}^n$  describe sequences of refining partitions of the interval  $[0, t]$ . The point  $\theta_k$  is chosen arbitrarily in  $[t_{k-1}, t_k]$ :

$$\theta_k = s t_k + (1 - s) t_{k-1} \quad \forall s \in [0, 1] \quad (2.4)$$

For ordinary Lebesgue-Stieltjes integrals the right hand side of (2.3) is independent of way  $\theta_k$  is sampled. In the present case, we instead have

- $\mathbb{L}^2(\Omega)$ -convergence of the sum

$$\begin{aligned} \sum_{t_k \in \mathfrak{p}} w_{\theta_k} (w_{t_k} - w_{t_{k-1}}) &= \sum_{t_k \in \mathfrak{p}} w_{\theta_k} (w_{t_k} - w_{\theta_k}) + \sum_{t_k \in \mathfrak{p}} w_{\theta_k} (w_{\theta_k} - w_{t_{k-1}}) \\ &= \sum_{t_k \in \mathfrak{p}} \left[ \frac{w_{t_k}^2 - w_{\theta_k}^2}{2} - \frac{(w_{t_k} - w_{\theta_k})^2}{2} \right] + \sum_{t_k \in \mathfrak{p}} \left[ \frac{w_{\theta_k}^2 - w_{t_{k-1}}^2}{2} + \frac{(w_{\theta_k} - w_{t_{k-1}})^2}{2} \right] \end{aligned} \quad (2.5)$$

whence it follows

$$\sum_{t_k \in \mathfrak{p}} w_{\theta_k} (w_{t_k} - w_{t_{k-1}}) = \frac{w_{t_n}^2}{2} - \sum_k \left[ \frac{(w_{t_k} - w_{\theta_k})^2}{2} - \frac{(w_{\theta_k} - w_{t_{k-1}})^2}{2} \right]$$

and in  $\mathbb{L}^2(\Omega)$  sense

$$\int_0^t w_s^{(\theta)} dw_s = \frac{w_t^2}{2} - \frac{t(1-2s)}{2}$$

- Average:

$$\begin{aligned} \mathbb{E} \int_0^t w_s^{(\theta)} dw_s &= \lim_{|\mathfrak{p}| \downarrow 0} \sum_{t_k \in \mathfrak{p}} \mathbb{E} \{ w_{\theta_k} w_{t_k} - w_{\theta_k} w_{t_{k-1}} \} \\ &= \sum_{t_k \in \mathfrak{p}} (\theta_k - t_{k-1}) = \sum_{t_k \in \mathfrak{p}} s(t_k - t_{k-1}) = \sigma^2 s \end{aligned}$$

### 3 Ito integral

Let suppose that  $\xi_t$  is a stochastic process satisfying the properties

1. **mean square integrability:**

$$\mathbb{E} \int_0^t ds \xi_s^2 < \infty$$

2. **Non anticipating:**  $\xi_t$  may depend only on  $w_s$  with  $s \leq t$ . As a consequence  $\xi_t$  and  $dw_t$  are **independent variables**

$$\mathbb{E} \xi_t dw_t = \mathbb{E} \xi_t \mathbb{E} dw_t = 0$$

**Definition 3.1.** For any stochastic process  $\xi_t$  satisfying the above two properties we can define the **Ito integral**

$$\int_0^t dw_s \xi_s := \lim_{|\mathfrak{p}| \downarrow 0} \sum_{t_k \in \mathfrak{p}} \xi_{t_{k-1}} (w_{t_k} - w_{t_{k-1}}) \quad (3.1)$$

Note that the approximating sums

$$I_n = \sum_{t_k \in \mathfrak{p}} \xi_{t_{k-1}} (w_{t_k} - w_{t_{k-1}}) \quad (3.2)$$

are defined in the Ito prescription by setting  $s$  to zero in (4.1). The convergence of (3.1) has to be understood in the **mean square sense** i.e.

$$\mathbb{E} (I_n - I_m)^2 \xrightarrow{n, m \uparrow \infty} 0$$

The definition (3.1) entails

i the *martingale* property

$$\mathbb{E} \int_0^t dw_s \xi_s = 0 \quad (3.3)$$

ii the *mean square integrability* property

$$\mathbb{E} \left( \int_0^t dw_s \xi_s \right)^2 = \int_0^t ds \mathbb{E} \xi_s^2 = \mathbb{E} \int_0^t ds \xi_s^2 \quad (3.4)$$

Namely

$$\mathbb{E} \left( \int_0^t dw_s \xi_s \right)^2 = \mathbb{E} \lim_{|\mathcal{p}| \downarrow 0} \sum_{t_k, t_l \in \mathcal{p}} (w_{t_{k+1}} - w_{t_k})(w_{t_{l+1}} - w_{t_l}) \xi_{t_k} \xi_{t_l} \quad (3.5)$$

As by hypothesis  $\xi_t$  is non anticipating

$$\mathbb{E} \left\{ (w_{t_{k+1}} - w_{t_k})(w_{t_{l+1}} - w_{t_l}) \xi_{t_k} \xi_{t_l} \right\} = \delta_{k,l} (t_{k+1} - t_k) \mathbb{E} \xi_{t_k}^2 \quad (3.6)$$

Hence upon inverting the limit and expectation value operation

$$\mathbb{E} \left( \int_0^t dw_s \xi_s \right)^2 = \lim_{|\mathcal{p}| \downarrow 0} \sum_{t_k, t_l \in \mathcal{p}} \delta_{k,l} (t_{k+1} - t_k) \mathbb{E} \xi_{t_k}^2 = \int_0^t dt \mathbb{E} \xi_t^2 \quad (3.7)$$

which yields the claim. Note that working directly in the continuum limit the above manipulations imply the rule

$$\mathbb{E} (dw_s dw_{s'}) = ds' \delta(s - s') \quad (3.8)$$

Furthermore since

$$\sum_{k=1}^n w_{t_{k-1}} (w_{t_k} - w_{t_{k-1}}) = \sum_{t_k \in \mathcal{p}} \frac{w_{t_k}^2 - w_{t_{k-1}}^2}{2} - \sum_{t_k \in \mathcal{p}} \frac{(w_{t_k} - w_{t_{k-1}})^2}{2}$$

in the *mean square sense* we can conclude

$$\int_0^t dw_s w_s = \frac{w_t^2}{2} - \frac{t}{2}$$

at variance with what expected from the ordinary rules of differential calculus. The origin of the discrepancy from ordinary calculus stems from

$$dw_t \sim O(\sqrt{dt})$$

**Example 3.1** (*Non-anticipative vs anticipative*). Let  $w_t$  a Wiener process for all  $t \geq 0$ , the function

$$f(t) = \begin{cases} 0 & \text{if } \max_{0 \leq s \leq t} w_s \leq 1 \\ 1 & \text{if } \max_{0 \leq s \leq t} w_s > 1 \end{cases}$$

is *non-anticipative* as it depends on the Wiener process up to the time  $t$  when the function is evaluated. On the other hand for any  $T > t$  the function

$$g(t) = \begin{cases} 0 & \text{if } \max_{0 \leq s \leq T} w_s \leq 1 \\ 1 & \text{if } \max_{0 \leq s \leq T} w_s > 1 \end{cases}$$

is *anticipative* as it depends on realizations of the Wiener process for times  $s$  posterior to the sampling time  $t$ .

**Example 3.2 (Exponential process).** Let us consider the process

$$\xi_t = e^{\lambda w_t - \frac{\lambda^2 t}{2}} \xi_o \quad (3.9)$$

by Ito lemma we have

$$d\xi_t = \lambda dw_t e^{\lambda w_t - \frac{\lambda^2 t}{2}} \xi_o = \lambda \xi_t dw_t$$

If we recast the Ito differential into Doob-Meyer form we find

$$\xi_t = \xi_o + \lambda \int_0^t dw_s \xi_s$$

The exponential process does not have bounded variation component.

## 4 The Stratonovich integral

We have seen that for

$$\theta_k = s t_k + (1 - s) t_{k-1} \quad \forall s \in [0, 1] \quad (4.1)$$

the sum

$$\sum_{t_k \in \mathcal{P}} w_{\theta_k} (w_{t_k} - w_{t_{k-1}}) = \frac{w_t^2}{2} - \sum_k \left[ \frac{(w_{t_k} - w_{\theta_k})^2}{2} - \frac{(w_{\theta_k} - w_{t_{k-1}})^2}{2} \right]$$

in  $\mathbb{L}^2(\Omega)$  converges to

$$\int_0^t w_s^{(\theta)} dw_s = \frac{w_t^2}{2} - \frac{t(1-2s)}{2}$$

Choosing  $s = 1/2$  the second term on the right hand side disappears and we recover the result from ordinary calculus. The example suggests to define the Fisk-Stratonovich integral

$$\int_0^t dw_s \diamond \xi_s := \lim_{|\mathcal{P}| \downarrow 0} \sum_{t_k \in \mathcal{P}} \xi_{\frac{t_{k-1} + t_k}{2}} (w_{t_k} - w_{t_{k-1}}) \quad (4.2)$$

Note that

$$\begin{aligned} & \xi_{\frac{t_{k+1} + t_k}{2}} - \frac{\xi_{t_{k+1}} + \xi_{t_k}}{2} = \\ & \frac{\xi_{t_{k+1} - \frac{t_{k+1} - t_k}{2}} - \xi_{t_{k+1}}}{2} + \frac{\xi_{t_k + \frac{t_{k+1} - t_k}{2}} - \xi_{t_k}}{2} = O(\xi_{t_{k+1}} - \xi_{t_k})^2 \end{aligned}$$

Thus we can equivalently write

$$\int_0^t dw_s \diamond \xi_s := \lim_{|\mathcal{P}| \downarrow 0} \sum_{t_k \in \mathcal{P}} \frac{\xi_{t_{k+1}} + \xi_{t_k}}{2} (w_{t_k} - w_{t_{k-1}})$$

As in the Ito case the limit converges in mean square sense. At variance with the Ito case, the integrand in the definition (4.2) is **anticipating**:

$$\mathbb{E} \xi_t \diamond dw_t \neq \mathbb{E} \xi_t \mathbb{E} dw_t = 0$$

Thus the martingale property of the Ito integral is lost. To appreciate the advantage of the definition consider

$$\int_0^t dw_s \diamond w_s = \lim_{|\mathcal{P}| \downarrow 0} \sum_{t_k \in \mathcal{P}} \frac{(w_{t_k} + w_{t_{k-1}})(w_{t_k} - w_{t_{k-1}})}{2} = \frac{w_t^2}{2} \quad (4.3)$$

in agreement with the rules of *ordinary differential calculus*. The example illustrates the general situation.

## 4.1 Relation with the Ito differential

Let us consider

$$\xi_t = g(\chi_t, t) \quad (4.4)$$

with

$$d\chi_t = b(\chi_t, t) dt + \sigma(\chi_t, t) dw_t \quad (4.5)$$

then by Ito lemma we can write

$$d\xi_t = dg(\chi_t, t) = dt \left\{ \partial_t + b_t \partial_{\chi_t} + \frac{\sigma^2}{2} \partial_{\chi_t}^2 \right\} g + dw_t \sigma_t \partial_{\chi_t} g \quad (4.6)$$

and use the this result to establish the relation between the Fisk-Stratonovich and the Ito integral. Namely given a non-anticipating process  $\eta_t$  we can couch the definition of the Fisk-Stratonovich integral into the form

$$\int_0^t dw_s \diamond \eta_s = \lim_{|\mathcal{p}| \downarrow 0} \sum_{t_k \in \mathcal{p}} \left\{ \eta_{t_{k-1}} (w_{t_k} - w_{t_{k-1}}) + \frac{(\eta_{t_{k-1}} - \eta_{t_k})(w_{t_k} - w_{t_{k-1}})}{2} \right\}$$

In the literature the latter equality is sometimes written in the continuum limit as

$$\int_0^t dw_s \diamond \eta_s = \int_0^t dw_s \eta_s + \langle \eta, w \rangle_t$$

where  $\langle \xi, w \rangle_t$  is *quadratic co-variation* of the processes  $\xi_t$  and  $w_t$ . The essential point is that in the limit (which converges in the mean square sense under our hypotheses) the quadratic co-variation receives finite contributions only from the term proportional to the increment of the Wiener process

$$dw_t \sim O(\sqrt{dt}) \quad \Rightarrow \quad dw_t^2 \sim O(dt) \quad (4.7)$$

In such a case if

$$\eta_t = f(\chi_t, t)$$

we find

$$\int_0^t dw_s \diamond f(\chi_s, s) = \int_0^t dw_s f(\chi_s, s) + \frac{1}{2} \int_0^t ds \sigma(\chi_s, s) \partial_{\chi_s} f(\chi_s, s) \quad (4.8)$$

In particular for

$$f(\chi_t, t) = \sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t)$$

we obtain

$$\begin{aligned} dw_t \diamond [\sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t)] &= \\ dw_t \sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t) &+ \frac{dt}{2} \sigma(\chi_t, t) \partial_{\chi_t} [\sigma(\chi_t, t) \partial_{\chi_t} g(\chi_t, t)] \end{aligned}$$

which allows us to write

$$d\xi_t = dg(\chi_t, t) = dt \left\{ \partial_t + \left[ b - \frac{\sigma}{2} (\partial_{\chi_t} \sigma) \right] \partial_{\chi_t} \right\} g + dw_t \diamond [\sigma \partial_{\chi_t} g] \quad (4.9)$$

As expected, the right hand side does not include any-longer a second derivative of  $g$ , the hallmark of Ito lemma. The function  $g$  is, however, transported by the Stratonovich stochastic differential

$$d\xi_t = dt \left[ b_t - \frac{1}{2} (\sigma \partial_{\chi_t} \sigma)_t \right] + dw_t \sigma_t$$

## 4.2 Examples

- Consider the process

$$\xi_t = w_t^2$$

In such a case the role of the process  $\chi_t$  of the previous section is played by the Brownian motion itself

$$\chi_t = w_t \quad \Rightarrow \quad d\chi_t = dw_t$$

i.e.  $b = 0$  and  $\sigma = 1$  in (4.5). Ito lemma yields

$$d\xi_t = dg(w_t) = 2 w_t dw_t + dt$$

The differential admits the **equivalent** Stratonovich representation

$$d\xi_t = 2 w_t \diamond dw_t$$

with again  $\chi_t = w_t$ .

- Consider now

$$\xi_t = \chi_t^2$$

with

$$d\chi_t = \chi_t dt + \chi_t dw_t \tag{4.10}$$

This case corresponds to

$$b = \sigma = \chi_t$$

in (4.5). It follows by Ito lemma

$$d\xi_t = 3\chi_t^2 dt + 2\chi_t^2 dw_t$$

On the other hand (4.10) admits the Stratonovich representation

$$d\chi_t = \frac{\chi_t}{2} dt + \chi_t \diamond dw_t$$

whence

$$d\xi_t = \chi_t^2 dt + 2\chi_t^2 \diamond dw_t \tag{4.11}$$

## References

- [1] L. C. Evans. An introduction to stochastic differential equations. UC Berkeley, Lecture Notes, 2006.
- [2] C. W. Gardiner. *Handbook of stochastic methods for physics, chemistry and the natural sciences*, volume 13 of *Springer series in synergetics*. Springer, 2 edition, 1994.